# Holographic Signatures of Resolved Cosmological Singularities: Numerical Investigations

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based on

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with N. Bodendorfer & J. Münch

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## Introduction: QG and Holography

Holographic aspects of QG actively investigated.

#### Main research directions (in LQG)

1. Dual statistical models from 3+0 QG partition function

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[Costantino '11; Dittrich, Hnybida '13; Bonzom, Costantino, Livine '15; Dittrich, Goeller,
Livine, Riello '17]
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2. Spin networks, tensor networks and holographic entanglement entropy

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[Han, Hung '16; Han, Huang '17; Chirco, Oriti, Zhang '17]
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3. Holography in symmetry-reduced models

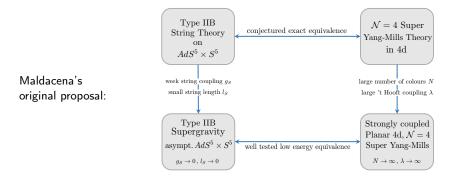
```
[Ashtekar, Wilson-Ewing '08; Bodendorfer, Schäfer, Schliemann '16; Bodendorfer, Mele, Münch '18]
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Related work: [Freidel '08; Bodendorfer, Thiemann, Thurn '11; Bodendorfer '15; Freidel, Perez, Pranzetti '16; Smolin '16; Livine '17, . . . ]

In this talk we will focus on the third direction.

# Gauge/Gravity Duality in a Nutshell

[Maldacena '97; Gubser, Klebanov, Polyakov '98; Witten '98]



## Holographic Principle ['t Hooft '93; Susskind '95; Bousso '02; ...]

Duality between gauge and (quantum) gravity theories

- ► Tools for QFT (e.g, condensed matter, QCD)
- ► Quantum gravity models

## QG for AdS/CFT

- Gauge/Gravity mainly understood in the classical (super)gravity regime
- Singularities in classical gravity
  - ▶ limitations in application of perturbative string theory
  - no consensus on the fate of singularities in dual QFT [Hertog, Horowitz '04, '05; Das, Michelson, Narayan, Trivedi '07; Barbón, Rabinovici '11; Craps, Hertog, Turok '12; Smolkin, Turok '12; Engelhardt, Horowitz '15, '16; ...]
- Progress in singularity resolution from non-perturbative QG

#### Question

QG resolution of bulk singularity  $\rightarrow$  signatures/improvements in dual QFT?

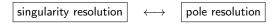
Prototype calculation: effective bulk quantum geometry in a cosmological setting

## Plan of the Talk

- 1. Setup and strategy
- 2. Holographic signatures of cosmological singularities

$$\boxed{ \text{bulk singularity} } \quad \longleftrightarrow \quad \boxed{ \text{pole in 2-point correlator} }$$

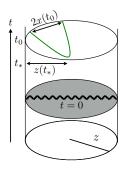
- 3. Resolved singularities in quantum corrected effective geometries
  - Simplifications needed for analytic treatment
  - ► Otherwise numerics necessary



4. Conclusions and outlook

## Setup

[Engelhardt, Horowitz '14; Engelhardt, Horowitz, Hertog '15]



Bulk Kasner-AdS: 
$$ds_5^2 = \frac{1}{z^2} (dz^2 + ds_4^2(t))$$

$$\textbf{Boundary:} \quad ds_4^2(t) = -dt^2 + \sum_{i=1}^3 \ t^{2p_i} dx_i^2 \quad , \quad p_i \in \mathbb{R}$$

Kasner conditions: 
$$\sum_{i} p_{i} = 1 = \sum_{i} p_{i}^{2}$$

 $p_{\it i} < 0$  geodesics bent towards the singularity

## Equal-time 2-point correlator

Geodesic approximation [Balasubramanian, Ross '00]

$$\langle \mathcal{O}(x)\mathcal{O}(-x)\rangle = \exp\left(-\Delta L_{ren}\right)$$

 $\Delta \equiv$  conformal weight of massive scalar operator  $\mathcal{O}$   $L_{ren} \equiv$  renormalized length of bulk geodesic connecting x and -x

# Renormalized Geodesic Length

fixed z metric

boundary metric

$$\frac{1}{z^2}\Big(-dt^2+a^2(t)\,dx^2+\ldots\Big)$$

$$dt^2 + a^2(t) dx^2 + \dots$$

Two-point correlator:

$$\langle \mathcal{O} \mathcal{O} \rangle_{CFT} = \mathbf{z}^{-2\Delta} \langle \phi \phi \rangle_{bulk}$$

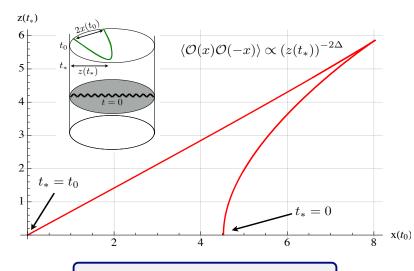
• Geodesic approximation:

$$\langle \phi \phi \rangle_{bulk} = \exp(-\Delta L)$$
 ,  $L \stackrel{\epsilon \to 0}{=} 2 \log(2z(t_*)) - 2 \log \epsilon$ 

$$\left\langle \mathcal{O}(-x)\mathcal{O}(x)
ight
angle_{\mathit{CFT}} = \exp\left(-\Delta\,L_{\mathit{ren}}
ight) = \left(2z(t_*)
ight)^{-2\Delta}$$

# Holographic Signature of Cosmological Singularity

[Engelhardt, Horowitz '14; Engelhardt, Horowitz, Hertog '15]



Finite distance pole in two-point correlator

## Quantum Corrected Kasner-AdS

#### Classical metric

$$ds_5^2 = \frac{1}{z^2} (dz^2 + ds_4^2(t))$$
 ,  $ds_4^2(t) = -dt^2 + \sum_{i=1}^3 t^{2p_i} dx_i^2$ 

- no large curvatures associated with z-direction
- curvature singularity in t-direction



#### Quantum corrected metric

$$ds_5^2 = \frac{1}{z^2} (dz^2 + ds_4^2(t, z))$$
 ,  $ds_4^2(t, z) = -dt^2 + a(t, z)^2 dx^2 + \dots$ 

- z-direction classical
- effective QG for 4d part s.t. singularity resolved

## Quantum Corrected Kasner-AdS

In principle effective metric from full 5d quantum Einstein eqs.  $\,\,
ightarrow\,$  here ansatz

- singularity resolution in  $ds_4^2 \longrightarrow \text{singularity-free } ds_5^2$  (possible z-dependence of  $ds_4^2$  must be small!)
- solution of 5d Einstein eqs. up to quantum corrections
- bulk QG effects relevant at 5d-Planck scale
- classical Kasner boundary metric
- quantum bounce with Kasner transitions
   [Gupt, Singh '12; Ashtekar, Wilson-Ewing '08; Wilson-Ewing '17]

## Quantum Corrected Metric: A First Example

#### Quantum corrected metric [Bodendorfer, Schäfer, Schliemann '16]

$$ds_5^2 = \frac{1}{z^2} \left( dz^2 + ds_4^2(t) \right) \quad , \quad a(t) = \frac{a_{\text{ext}}}{\lambda^p} \left( t^2 + \lambda^2 \right)^{p/2}$$

Two main simplifications:

- 1.  $\lambda$  related to 4d bulk Planck scale
- 2. No Kasner transitions

#### Renormalized length and two-point correlator

Geodesic eqs. analytically solvable:

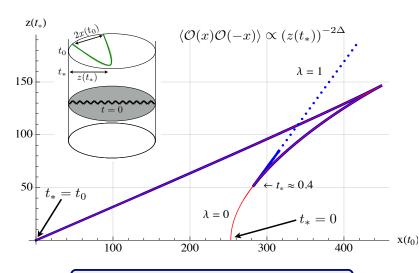
$$\pm s(z=\epsilon) \stackrel{\epsilon \to 0}{=} \log{(2z(t_*))} - \log{(\epsilon)}$$
  $\epsilon = \text{boundary regulator}$ 

and

$$L_{ren} = 2 \log (2z(t_*))$$
  $\Rightarrow$   $\langle \mathcal{O}(x)\mathcal{O}(-x)\rangle = (2z(t_*))^{-2\Delta}$ 

# Holographic Signature of Resolved Cosmological Singularity

[Bodendorfer, Schäfer, Schliemann '16]



No finite distance pole in two-point correlator!

## **Drop Simplification 1: 5d Planck Scale**

#### Quantum corrected metric [Bodendorfer, Mele, Münch '18]

$$ds_5^2 = \frac{1}{z^2} \left( dz^2 - dt^2 + \sum_i \frac{a_i^2(t,z)}{a_i^2(t,z)} dx_i^2 \right) \quad , \quad a_i^2(t,z) = \frac{a_{ext}^2}{\lambda^2 p_i} \left( t^2 + \frac{z^2}{\lambda^2} \lambda^2 \right)^{p_i}$$

•  $ds_5^2$  not singular at t=0

$$R_{\rm Kretschmann}^{(5)} = z^4 R_{\rm Kretschmann}^{(4)} + \dots \quad , \quad R_{\rm Kretschmann}^{(4)} \sim \lambda^{-4}$$

- $\Rightarrow$  effective 5d scale:  $\lambda \to z\lambda$  s.t.  $R_{\text{Kretschmann}}^{(5)} = \mathcal{O}(\lambda_{5d}^{-4})$
- z-dependence of  $ds_4^2$  small (order of  $\lambda$ ) and  $\partial_z a \to 0$  in zeroth order in  $\lambda$
- classical Kasner-AdS boundary metric recovered  $(\frac{a_{\rm ext}}{\lambda p} o 1$  ,  $\lambda, z o 0)$
- symmetries:
  - ► translation (e.g., for us in x-direction)
  - ▶ scaling:  $t \longmapsto \Lambda t$  ,  $z \longmapsto \Lambda z$  ,  $x \longmapsto \Lambda^{1-p} x$   $(\Lambda \in \mathbb{R})$
- no Kasner transitions

# **Boundary Pole Resolution**

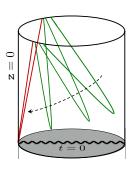
Boundary metric is singular

$$a_i^2(t,z) = rac{a_{\mathrm{ext}}^2}{\lambda^{2p_i}} \left(t^2 + \lambda^2 z^2\right)^{p_i} \stackrel{z o 0}{\longrightarrow} rac{a_{\mathrm{ext}}^2}{\lambda^{2p_i}} t^{2p_i}$$

Classical Kasner-AdS: [Engelhardt, Horowitz, Hertog '15]

$$z_* \rightarrow 0$$
 as  $t_* \rightarrow 0$ 

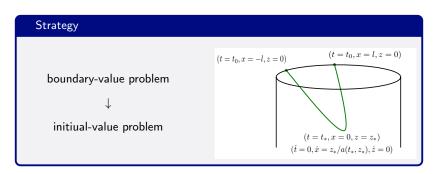
family of space-like geodesics approaches a light-like boundary geodesic probing high-curvature regime



Quantum corrected bulk: **boundary geodesics isolated** (more later).

# **Solving Geodesic Equations**

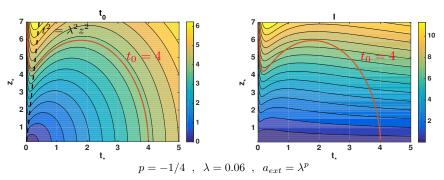
geodesic Eqs. not decoupled  $\longrightarrow$  numerics necessary!



- ullet initial values at the turning point only in terms of  $t_*$  and  $z_*$
- ullet solve IVP numerically and find  $t_0(t_*,z_*)/I(t_*,z_*)$
- we are interested in fixed  $t_0$ -values and vary  $l o find relation <math>z_*(l)$

# **Solving Geodesic Equations**

**Key point:** relate turning point data  $(t_*, z_*)$  with boundary data  $(t_0, I)$ 

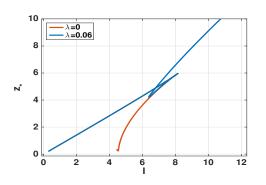


#### Note

- Classical case: bulk geodesics approaching a null boundary geodesic  $(z_* \to 0 \text{ for } t_* \to 0)$
- Quantum case:  $z_*$  grows as  $t_* \to 0 \Rightarrow$  null boundary geodesic not limit of bulk geodesics

# **Solving Geodesic Equations**

Reading off the values of l along  $t_0$ -level line, we get the relation  $z_*(l)$ :



$$ho = -1/4$$
 ,  $a_{\rm ext}/\lambda^p = 1$   $\lambda = 0$   $\lambda = 0.06$ 

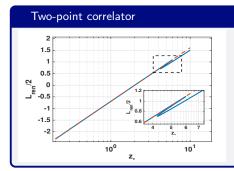
last step: check 
$$L_{ren} \sim \log \left( z(t_*) \right) \Rightarrow \langle \mathcal{O}(-x) \mathcal{O}(x) \rangle \propto \left( z(t_*) \right)^{-2\Delta}$$

# Renormalized Geodesic Length & Two-point Correlator

## Renormalization of geodesic length

- ullet Conformal boundary (z=0) lies at infinity o geodesic length diverges
- Renormalization procedure needed:
  - 1. Evaluate geodesic up to  $\epsilon = z_{UV}/t << 1$
  - 2. Subtract divergent contribution ( $\epsilon \to 0$ )

$$L_{ren} = L - L_0$$
 ,  $L_0 = 2 \log \left( \left| \frac{z_{UV}}{t} \right| \right)$ 



$$\langle \mathcal{O}(-I)\mathcal{O}(I) 
angle \propto e^{-\Delta L_{ren}} = (z_*)^{-2\Delta} + z_* 
eq 0$$
 for finite non-zero  $I$ 

pole resolved

## **Drop Simplification 2: Inclusion of Kasner Transitions**

Inclusion of Kasner transitions highly non trivial:

- Find a metric satisfying all requirements discussed before not straightforward
- We provided two examples of quantum corrected metrics partially satisfying them

 $\downarrow$ 

pole resolution qualitatively the same!

## Conclusion

## Summary

- $\bullet$  Bulk singularity resolution  $\stackrel{\textit{dual}}{\longleftrightarrow}$  boundary finite-distance pole resolution
- Application of LQG in Gauge/Gravity-framework possible
- New tools for holography/improvements in dual field theory

#### Conclusion

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#### Future work

- 5d-Loop quantization
- Look for independent field theory calculations, e.g. via lattice QFT
- Generalize to other spacetimes, e.g. black holes (work in progress)
- Long term goal: finite N gauge theory

#### Conclusion

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## Thank you for your attention!

## **Inclusion of Kasner Transitions**

#### Quantum corrected metric

$$ds_5^2 = \frac{1}{z^2} \left( dz^2 - dt^2 + a(t, z)^2 dx^2 + \dots \right)$$

with

$$\mathit{a}(t,z)^2 = \frac{\mathit{a}_{\mathrm{ext}}^2}{\lambda^{2p}} \left( t^2 + \lambda^2 z^2 \right)^p \exp \left[ 2 \Delta p \sinh^{-1} \left( \frac{t}{z \lambda} \right) \right] \ , \quad \Delta p \in \mathbb{R}$$

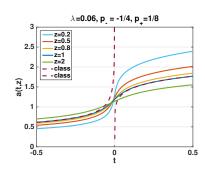
- $ds_4^2$ :  $a(t,z) = \frac{a_{ext}}{\lambda^p} \left(t^2 + \lambda^2\right)^{p/2} \exp\left[2\Delta p \sinh^{-1}\left(\frac{t}{\lambda}\right)\right]$ 
  - limiting curvature mimetic gravity [Chamseddine, Mukhanov '16]
  - toy model effective homog. and isotr. LQG [Bodendorfer, Schäfer, Schliemann '17]
- $\lambda \to z\lambda$  (quantum effects at 5d Planck scale)
- implement transition  $p_-=p-\Delta p \ 
  ightarrow \ p_+=p+\Delta p$
- $\Delta p = 0 \, o \, 5 d$  Planck scale quantum corrected metric with no transitions

## Inclusion of Kasner Transitions

For  $\left|\frac{t}{\tau\lambda}\right| >> 1$ :

$$a(t,z)^2 \simeq rac{a_{ext}^2}{\lambda^{2p}} igg(rac{2}{\lambda z}igg)^{\pm 2\Delta p} t^{2p\pm} \quad , \quad p_\pm = p \pm \Delta p$$

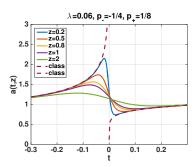
 $\Rightarrow$  transition  $p_- \to p_+$  at the bounce (t=0).



boundary 4d classical Kasner	X
solution of 5d Einstein eqs. up to $\mathcal{O}(\lambda)$	×
Kretschmann condition $R_{Kretsch.}^{(5)} = \mathcal{O}(\lambda_{5d}^{-4})$	<b>✓</b>
Qualitative LQC behavior	<b>✓</b>
symmetries (translation, scaling)	<b>✓</b>

## **Alternative Metric for Kasner Transitions**

$$a(t,z)^{2} = \frac{a_{\text{ext}}^{2}}{\lambda^{2p}} \left( t^{2} + \lambda^{2} z^{2} \right)^{p} \exp \left[ 2\Delta p \sinh^{-1} \left( \frac{t}{z\lambda} \right) \right] \cdot \underbrace{\frac{2}{8}}_{\text{ex}} 1.5$$
$$\cdot (\lambda z)^{2\Delta p \tanh \left( \frac{t}{z\lambda} \right)}$$

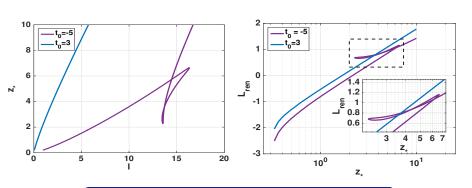


boundary 4d classical Kasner	<b>✓</b>
solution of 5d Einstein eqs. up to $\mathcal{O}(\lambda)$	<b>✓</b>
qualitative LQC behavior	×
Kretschmann condition $R_{Kretsch.}^{(5)} = \mathcal{O}(\lambda_{5d}^{-4})$	×

#### Numerical Results with Kasner Transitions

#### Two kinds of solutions:

- starting at  $t_0 < 0$ , bent towards the resolved singularity and passing it
- starting at  $t_0 > 0$ , bent away from the resolved singularity



Finite-distance pole in two-point correlator resolved!

## **Numerical Checks**

## cut-off dependence?

#### No Kasner Transitions

Z <sub>UV</sub>	$\exp\left(-L_{ren}\right)$	$\exp\left(-L_{ren}\right)$
0.05	0.761810	0.917559
0.06	0.761841	0.917597
0.07	0.761860	0.917661
0.08	0.761896	0.917716
0.09	0.761943	0.917783
0.1	0.761972	0.917897

$$\lambda = 0.06, p = -\frac{1}{4}$$

#### Kasner Transitions

$\exp\left(-L_{ren}\right)$	$\exp\left(-L_{ren}\right)$
0.367113	0.308506
0.367149	0.308554
0.367185	0.308602
0.367229	0.308661
0.367267	0.308714
0.367314	0.308774
	0.367113 0.367149 0.367185 0.367229 0.367267

$$\lambda = 0.06 \; , \; p = -\frac{1}{16} \; , \; \Delta p = \frac{3}{16}$$

Renormalization procedure independent of  $z_{UV}$  within 0.1%

# Comparison with Analytical Results: No z-dependence in ds<sub>4</sub><sup>2</sup>

